The dark matter environment of the Abell 901/902 supercluster: a weak lensing analysis of the HST STAGES survey

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ABSTRACT

We present a high resolution dark matter reconstruction of the $z = 0.16$ Abell 901/902 supercluster from a weak lensing analysis of the STAGES HST survey. We detect the four main structures of the supercluster at high significance, resolving substructure within and between the clusters. We find that the distribution of dark matter is well traced by the cluster galaxies, with the brightest cluster galaxies marking out the strongest peaks in the dark matter distribution. We also find a significant extension of the dark matter distribution of Abell 901a in the direction of an infalling X-ray group Abell 901a. We present mass, mass-to-light and mass-to-stellar mass ratio measurements of the structures and substructures that we detect. We find no evidence for variation of the mass-to-light and mass-to-stellar mass ratio between the different clusters. We compare our space-based lensing analysis with an earlier ground-based lensing analysis of the supercluster to demonstrate the importance of space-based imaging for future weak lensing dark matter ‘observations’.

1 INTRODUCTION

Observations and theory both point to the importance of environment on the properties of galaxies. Early-type galaxies are typically found in more dense regions compared to late-type galaxies (Dressler 1980), galaxy colour and luminosity are found to be closely related to galaxy density (Blanton et al. 2005) and the fraction of star-forming galaxies also shows a strong sensitivity to the density on small $\sim 1$ Mpc scales (Balogh et al. 2004; Blanton et al. 2006). Theoretically there are a number of physical mechanisms that could cause these effects in dense environments. These processes can change the star formation history, gas content and/or morphology of a galaxy through, for example, ram-pressure stripping (Gunn & Gott 1972) and/or the tidal effects of nearby galaxies (galaxy harassment, Moore et al. 1996) and/or the tidal effects of the dark matter potential (Bekki 1999; Moore et al. 1998). They depend differently on cluster gas, galaxy density and the dark matter potential however, with ram-pressure stripping dependent on the gas distribution compared to tidal effects which are dependent...
on the overall potential. A key difficulty in disentangling these effects observationally, is that typically the tidal potential is only constrained in a global sense through the measured velocity dispersion of a cluster, or a richness or total luminosity estimate. This results in an assumed spherical tidal potential model that is smoothed over the small scales that are relevant for tidal stripping and harassment studies.

In this paper we study the complex Abell 901/902 supercluster, hereafter A901/2, in the first of a series of papers from the STAGES collaboration. From a rich multi-wavelength dataset the A901/2 supercluster permits a thorough investigation of the relationships between galaxy morphology (from Hubble Space Telescope and ground-based imaging, Gray et al. 2007b; Lane et al. 2007), luminosity, stellar mass and colour (from the COMBO-17 survey with 17-band optical imaging, Wolf et al. 2003; Borch et al. 2006), star formation rates (from 24\mu m Spitzer data Bell et al. 2007), galaxy density (Wolf et al. 2005; Gray et al. 2004), and the hot intra-cluster medium (from XMM observations, Gilmour et al. 2007; Gray et al. 2007a). One of the key reasons to obtain Hubble Space Telescope (HST) imaging of this supercluster was to construct a high resolution, reliable and accurate map of the projected total mass density distribution. Using weak gravitational lensing techniques we are able to reconstruct the distribution of both dark and luminous matter and quantify the significance of the structures that are seen, updating the previous ground-based weak lensing analysis of Gray et al. (2002). This extra dimension to the multi-wavelength view of A901/2 will be a key ingredient in future studies where we hope to be able to separate the effects of tidal and gas-dynamical influence on galaxy formation and evolution.

Weak gravitational lensing is now a well-established method for studying the distribution of dark matter. Light from distant galaxies is deflected by the gravitational effect of the intervening structures, inducing a weakly coherent distortion in the shapes of galaxy images. The strength of this lensing effect is directly related to the projected mass along the line of sight, and it can therefore be used to map dark matter in dense regions (see for example Gray et al. 2002; Gavazzi et al. 2004; Dietrich et al. 2005; Clowe et al. 2006; Mahdavi et al. 2007).

The first weak lensing analysis of A901/2 by Gray et al. (2002) used deep ground-based R-band observations from the COMBO-17 survey (Wolf et al. 2003). This analysis revealed three significant peaks in the dark matter distribution at the locations of the A901a, A901b and A902 clusters, in addition to a low significance south west peak coincident with a galaxy group, hereafter referred to as the SW group. This analysis also showed a filamentary extension between the A901a and A901b clusters. As this filament was located across the CCD chip boundary in the mosaic image, however, Gray et al. (2002) could not rule out the possibility of this structure originating from residual uncorrected distortions from the point spread function (PSF) of the telescope and detector. Using the accurate photometric redshift information from the A901/2 17-band observations of the COMBO-17 survey, Taylor et al. (2004) extended the Gray et al. (2002) analysis, by creating a three-dimensional reconstruction of the A901/2 dark matter distribution. This analysis revealed a previously unknown higher redshift cluster located behind A902 that is at $z = 0.16$. This cluster was named and hereafter referred to as CB1 by Taylor et al. (2004), and we have updated the redshift based on an improved photometric redshift catalogue.

In this analysis we revisit the dark matter distribution in A901/2 using deep Hubble Space Telescope (HST) observations. The dominant source of noise in weak lensing cluster analysis is the shot noise introduced from the random intrinsic ellipticities of galaxies. Weak lensing maps of dark matter on small scales therefore benefit greatly from the high resolution that HST has to offer, as it boosts the number density of resolved galaxies from which the lensing signal can be measured, reducing the intrinsic ellipticity noise on small scales. In addition, the high resolution space-based data permits higher signal-to-noise shape measurements and a narrower PSF, thus implying a more accurate PSF correction.

This paper is organised as follows. In section 2 we describe the weak lensing theory that is related to this analysis, and the maximum likelihood method that we use to reconstruct the dark matter distribution. We describe the data and weak lensing measurement method in section 3. We present our results in section 4, including NFW profile mass measurements in section 4.1 and the dark matter reconstruction and ground-based comparison in section 4.2. A first comparison of the dark matter and galaxy light distribution is presented in section 4.3 along with mass, mass-to-light and mass-to-stellar mass ratio measurements. A more detailed comparison of the mass, gas and galaxies of A901/2 will appear in a forthcoming analysis. We investigate the significance of the supercluster substructure that is resolved in our dark matter reconstruction in section 4.4 and discuss our findings and conclude in section 5. Throughout this paper we assume a $\Lambda$CDM cosmology with $\Omega_m = 0.3$, $\Omega_{\Lambda} = 0.7$, and $H_0 = 100h$ km s$^{-1}$ Mpc$^{-1}$. All magnitudes are given in the Vega system.

2 METHOD AND THEORY

Gravitational lensing is sensitive to the projected surface mass density along the line of sight $\Sigma(\theta)$, typically denoted by the convergence $\kappa$. In the case of a single lens,

$$\kappa = \frac{\Sigma}{\Sigma_{\text{crit}}}, \quad \Sigma_{\text{crit}} = \frac{c^2 D_l}{4\pi G D_s D_{ts}},$$

where $D_l$ is the angular diameter distance to the lens, $D_s$ is the angular diameter distance to the lensed source galaxies, and $D_{ts}$ is the angular diameter distance from the lens to the source.

The coherent distortion, or reduced shear $g = g_1 + ig_2$, that is detected in the images of distant sources allows for the reconstruction of the projected intervening matter $\kappa$ as $g = \gamma(1 - \kappa)$, and

$$\kappa = \frac{1}{2}(\psi_{11} + \psi_{22}), \quad \gamma_1 = \frac{1}{2}(\psi_{11} - \psi_{22}), \quad \gamma_2 = \psi_{12},$$

where $\gamma$ is the true shear, $\gamma = \gamma_1 + i\gamma_2$, and $\psi_{ij}$ is the second derivative of the lensing potential (see for example Bartelmann & Schneider 2001).

The strength of all lensing distortions is invariant under the transformation $\kappa' = (1-\lambda)\kappa + \lambda\kappa_{0}$, where $\lambda$ is a constant (Gorenstein et al. 1988). This is known as the ‘mass sheet degeneracy’ implying that all lensing observations are insensitive to a constant mass sheet across the field of view. For wide-field images of relatively isolated clusters, one can remove this bias using the $\zeta_c$ statistic of Clowe et al. (1998). The $\zeta_c(r)$ statistic gives a model free estimate of the mass enclosed within a radius $r$ and is given by

$$\zeta_c(r_1) = \bar{\kappa}(r \leq r_1) - \bar{\kappa}(r_2 \leq r \leq r_{\text{max}}),$$

where $r_2$ is defined to be the radius outside which the cluster density is expected to be very low based on initial mass estimates.

1 ‘Space Telescope A901/902 Galaxy Evolution Survey’ (HST GO-10395, PI M. E. Gray), www.nottingham.ac.uk/~ppzmeg/stages/
and \( r_{\text{max}} \) is the field-of-view radius. The second term therefore essentially measures the constant \( \lambda \). In the case of A901/2 we find \( \kappa \left( 15' \leq r \leq 20' \right) = -0.002 \pm 0.007 \) where \( r \) is measured from the centre of the STAGES mosaic which is centred on the supercluster. This measure is consistent with what would be expected from large-scale structure and the NFW multi-halo model of the A901/2 supercluster that we develop in section 4.1. We therefore assume a zero mass sheet degeneracy correction in the analysis that follows.

2.1 Dark Matter reconstruction

In this paper we use a maximum likelihood method to reconstruct the surface mass density \( \kappa \). Starting with a ‘best guess’ Kaiser \& Squires (1993) reconstruction, the lensing potential \( \psi \) is constructed on a pixelised grid and is allowed to vary to produce the minimum difference between the reconstructed and observed reduced shear field. The benefit of using this method is that a varying noise estimate can be obtained across the whole region enabling the significance of each structure in the dark matter map to be accurately quantified. Furthermore it does not rely on the assumption that the observed reduced shear \( g \) is approximately equal to the true shear \( \gamma \), which for the A901/2 supercluster would introduce errors at the \( \sim 15\% \) level. We smooth the resulting \( \kappa \) maps with a Gaussian of smoothing scale 0.75 arcmin, which is equal to \( \sim 130 \) kpc at the supercluster redshift \( z = 0.16 \) (assuming a \( \Lambda \)CDM cosmology with \( h = 0.7 \)). This smoothing scale provides the best trade-off between high resolution and high signal-to-noise. A wavelet reconstruction (Starck et al. 2006) would have automatically optimised the smoothing scale across the field of view. As we are particularly interested in the significance of the structures revealed in the dark matter reconstruction, however, the wavelet method, which does not conserve noise properties, cannot be used.

Once peaks are detected in a weak lensing mass map their significance has to be compared to what is expected from a smoothed random noise map, with a 3\( \sigma \) noise peak, for example, is much more common than would naively be expected. As shown by Van Waerbeke (2000), the statistics of peaks in a smoothed pure noise map follow the peak statistics of a two-dimensional Gaussian random field (Bond \& Efstathiou 1987). We use both the peak signal-to-noise and the radial peak profile to calculate the global probability of a detected dark matter peak arising from noise using Equation (45) of Van Waerbeke (2000).

2.2 Model-free Mass measurement

As our dark matter reconstruction reveals structures that are far from the spherically symmetric simple isothermal sphere and NFW models (Navarro et al. 1997) that are often fit to estimate masses from weak lensing measurements (see for example Hoekstra 2007), our preferred method to measure mass uses a model-free mass estimate. Following the idea of the \( \zeta_c \) statistic (Equation 3), we measure the mass of structures within an aperture. For the main structures in the supercluster we define apertures by the \( \sigma \) and 3\( \sigma \) enclosed regions in the dark matter signal-to-noise maps. In the cases of smaller cluster substructure, where the smoothed structures appear to be more spherical, we use circular apertures of radius 0.75 arcmin to match the smoothing scale used in the dark matter reconstruction. The ‘aperture’ mass is given by

\[
M = \sum_{\text{aperture}} A_{\text{pix}} \kappa(x, y) \Sigma_{\text{crit}},
\]

where \( A_{\text{pix}} \) is the projected pixel area at the cluster redshift in \( h^{-2} \text{Mpc}^2 \), \( (x, y) \) are pixels enclosed by the chosen aperture and \( \Sigma_{\text{crit}} \) is the critical surface mass density, given in Equation 1.

2.3 NFW profile model

The main drawback of using the model-free mass estimate in Equation 4 is the inability to separate mass at different redshifts. This is because the dark matter reconstruction \( \kappa \) projects the measured surface mass density along the line of sight. In the case of the A901/2 supercluster there is a higher redshift \( z = 0.46 \) cluster, CB1, that lies behind A902 (Taylor et al. 2004) such that our model free mass estimate for the A902 region therefore gives us the combined mass of A902 and CB1. To obtain separate mass estimates for the A902 and CB1 cluster and to enable a comparison to future analyses of numerical simulations, we therefore also present mass estimates for the dark matter structures in A901/2 using an NFW halo model.

The NFW halo model has been shown in numerical simulations to provide a good fit to the spherically averaged profile of all dark matter halos irrespective of their mass (Navarro et al. 1997). The NFW model for the density profile of a halo at redshift \( z \) is given by

\[
\rho(r) = \frac{\delta_c \rho_c(z)}{(r/r_s)(1 + r/r_s)^2},
\]

where \( \delta_c \) is the characteristic density, \( r_s \) is the scale radius and \( \rho_c(z) \) is the critical density\(^2\) given by 3H(\(z\))^2/8\(\pi\)G. The virial radius \( r_\Delta \) is defined as the radius where the mass density of the halo is equal to \( \Delta \rho_c \), such that the corresponding virial mass \( M_{\text{NFW}} \) is given by

\[
M_{\text{NFW}} = \Delta \rho_c \frac{4\pi}{3} r_\Delta^3.
\]

Eke et al. (1996) show that for flat \( \Lambda \) cosmology, assuming a spherical collapse model,

\[
\Delta = 178 \Omega_m(z)^{0.45}
\]

For the cosmology adopted here \( \Delta = 118 \) at the A901/2 supercluster redshift \( z = 0.16 \), and \( \Delta = 138 \) at the CB1 cluster redshift \( z = 0.46 \). As the mass enclosed within a radius \( R \) is given by

\[
M(r \leq R) = 4\pi \delta_c \rho_c(z) r_s^3 \ln \left( 1 + \frac{R}{r_s} \right) - \frac{R}{r_s} \frac{M_{\text{vir}}}{1 + R/r_s}
\]

defining the concentration parameter as \( c = r_\Delta/r_s \), the characteristic halo density \( \delta_c \) is given by

\[
\delta_c = \frac{\Delta}{3} \ln (1 + c) - c/(1 + c)
\]

For a given CDM cosmology, the halo mass \( M_{\text{NFW}} \) and concentration \( c \) are related (Navarro et al. 1997; Bullock et al. 2001; Eke et al. 2001), where the dependence is calculated through fits to numerical simulations. In this paper we use the relationship between halo mass \( M_{\text{NFW}} \) and concentration \( c \) derived by Bullock et al. (2001).

The expression for the weak lensing shear \( \gamma \) and convergence \( \kappa \) induced by an NFW dark matter halo, given in Wright \& Brainerd (2000), depends on the redshift of both the lens and source galaxies. In this analysis we have accurate redshifts for each cluster member but no redshift information for \( \sim 90\% \) of our source galaxies.

\(^2\) Note that within the literature \( \rho_c \) is often defined differently. In this paper we follow the definition given in the Appendix of Navarro et al. (1997).
as they are too faint to calculate a COMBO-17 photometric redshift. The maximum-likelihood method of Schneider & Rix (1997) was designed to take advantage of such a data set for analysing the galaxy-galaxy lensing statistically (Kleinheinrich et al. 2006; Heymans et al. 2006), and it is this method that we have adapted for cluster lensing and describe below.

For a model cluster density profile, in the case where all galaxy redshifts are known, the weak shear $\gamma$ and convergence $\kappa$ experienced by each source galaxy can be predicted by summing up the shear and convergence contributions from all the foreground clusters. In this analysis the redshifts of the source galaxies are unknown, and we therefore assign those galaxies a magnitude-dependent redshift probability distribution $p(z; \text{mag})$ given by Equation 15 of Heymans et al. (2005) updated with the magnitude-redshift relation of Schrabback et al. (2007), where

$$z_m = 0.29|\alpha (606) - 22| + 0.31. \quad (10)$$

We are then able to calculate the expectation value of the observed reduced shear $\langle g \rangle$ through Monte Carlo integration by drawing a source galaxy redshift estimate $z^*_\nu$ from the distribution $p(z; \text{mag})$, $\nu = 1..N_{\text{MC}}$ times, where $N_{\text{MC}} = 100$ in this analysis. Testing larger values for $N_{\text{MC}}$ did not change the result. For each $z^*_\nu$ estimate the induced cluster lensing shear $g^{\nu}$ is calculated with the resulting mean reduced shear given by

$$\langle g \rangle = \frac{1}{N_{\text{MC}}} \sum_{\nu=1}^{N_{\text{MC}}} g^{\nu}. \quad (11)$$

The intrinsic source galaxy ellipticity $\epsilon^i$ is then calculated, $\epsilon^i \approx \epsilon^{\text{obs}} - g$. The distribution of each component of the observed galaxy ellipticity is well described, for the STAGES survey, by a Gaussian with errors $\epsilon^{\text{obs}}$ and $\epsilon^{\text{MC}} = \epsilon^{\text{MC}} - 0.03$ arcsecond pixel scale, as suggested by Rhodes et al. (2007). The best-fit dark matter parameters are determined by maximising the likelihood $L = \Pi \left[ P(\epsilon^i) \right]_i$ where the product extends over all source galaxies $i$.

## 3 THE STAGES DATA

The STAGES survey (Gray et al. 2007b) spans a quarter square degree centered on the A901/2 supercluster. Imaged in F606W, using the HST Advanced Camera for Surveys (ACS), the 80 orbit mosaic of 80 ACS tiles forms the second largest deep image taken by HST. A detailed account of the STAGES reduction method will be presented in Gray et al. (2007b). It is very similar to the reduction used for the GEMS survey discussed in Heymans et al. (2005) and Caldwell et al. (2007), differing only in the dither and drizzle strategy. For STAGES, each image consists of four co-added dithered images combined with a Gaussian drizzling kernel with a resulting 0.03 arcsecond pixel scale, as suggested by Rhodes et al. (2007). STAGES is complemented by 17-band imaging from the COMBO-17 survey which provides accurate photometric redshifts with errors $\sigma_z \sim 0.02(1 + z)$, spectral energy distribution galaxy classification, and stellar mass estimates $M_\star$ from low resolution 17-band spectra fits to parameterised star formation history models (Borch et al. 2006; Wolf et al. 2004).

### 3.1 Weak lensing shear measurement

To measure the reduced weak lensing shear $g$, we use the data reduction steps and method described in Heymans et al. (2005). The shear measurement aspect is based on the Kaiser et al. (1995) method. As we are primarily interested in the variation in the dark matter map we have updated our shear measurement pipeline to maximise the signal-to-noise by including a polynomial fit to the shear seeing correction $P^\gamma$ (Equation 7 in Heymans et al. 2006) as a function of galaxy size. We also include the Hoekstra correction to the shear polarisability tensor detailed in Heymans et al. (2006). The accuracy of these updates has been verified using the publicly available suite of simulations from the Shear TEsting Programme3 (Heymans et al. 2006; Massey et al. 2007). The modifications successfully reduced the noise on the shear measurement, quantified through the root-mean-square variation of the measured ellipticity $\sigma$, from $\sigma = 0.31$ to $\sigma = 0.26$.

We use the same method as Heymans et al. (2005) to account for the time variation of the ACS PSF, namely to divide the data into sets imaged in a short period of time and assume that the temporal variation during that time is minimal. The majority of the A901/2 field was observed in the space of 20 days, with the remaining 10% imaged at a later date over the space of 4 days. Owing to the relatively low galactic latitude of the A901/2 field and the resulting high stellar density of $30 - 40$ useful stellar images per ACS image, we are able to split the data into seven groups to achieve good temporal sampling of the PSF distortion. This number was chosen to balance between the need to use as many ACS images as possible to maximise the signal-to-noise on the average measured stellar ellipticity as a function of CCD position, whilst requiring as many time bins as possible to minimise the temporal variation of the PSF.

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pattern. Figure 1 shows the tiling pattern of the STAGES ACS observations denoting each group of data that was used to make the seven different PSF models. With this semi-time dependent model we find and remove temporal variation during the A901/2 observations. Averaged across the ACS field-of-view, this temporal variation is at the 1% level on the measured stellar ellipticity. As this variation is more than an order of magnitude lower than the weak lensing signal from the A901/2 supercluster our semi-time dependent PSF model is more than sufficient for this analysis. We might expect to see low-level systematics for the ACS images whose observation date is isolated at the start or end of a data group, affecting tiles 21, 33, 36, 43, 44, 46, 47, 57, 69 and 72. Indeed in the analysis that follows we find \( > 3.5\sigma \) B-modes, an indication of systematics (Crittenden et al. 2002), in tiles 33, 36, 44, 46, 57. As these tiles do not co-incide with the four main supercluster structures, however, they do not effect the majority of our analysis which we discuss further in section 4. We also include a conservative systematic error term based on the B-mode amplitude in the analysis that follows. Note that Schrabback et al. (2007) and Rhodes et al. (2007) present significantly more advanced methods to model the temporal variation of the ACS PSF designed for the detection of the weaker lensing signal from large-scale structure.

In the time since the ACS observations used by Heymans et al. (2005) were imaged, the charge transfer efficiency (CTE) of the ACS has degraded significantly. During the CCD readout, as the charge transfer is independent of the pixel count, CTE impacts on the readout amplifiers. As the amount of charge left behind in each charge transfer is independent of the pixel count, CTE impacts on the readout amplifiers and SN is a signal-to-noise estimate that we define as the ratio of the flux and flux error measurements from SEXtractor (Bertin & Arnouts 1996). As a normalisation constant derived to minimise the average measured shear \((g_1)\), where \(g_1 = 2(e_{1} - e_{1}^{\text{CTE}}) / \text{Tr}(P^T)\), \(e_1\) is the PSF corrected galaxy ellipticity and \(P^T\) is the shear polarisability tensor from Luppino & Kaiser (1997). For galaxies furthest from the readout amplifier, and hence the most strongly affected, this is a percent level correction to the measured ellipticity, which is more than an order of magnitude lower than the weak lensing signal from the A901/2 supercluster.

3.2 Galaxy selection and redshift estimation

As we are interested in the dark matter in A901/2 at a redshift of \(z = 0.16\), we select galaxies that are likely to be at higher redshifts and thus lensed by the supercluster. As the majority of our galaxies are too faint to calculate a COMBO-17 photometric redshift, the best option is to use a magnitude selection. We use the relationship between median redshift \(z_m\) and F606W magnitude derived in Schrabback et al. (2007) and given in Equation 10. To ensure that the majority of objects have \(z_s > z_{A901/2}\), we select galaxies with \(m_{\text{F606}} > 23\), corresponding to a median redshift \(z_m > 0.6\). We also include selection criteria chosen to optimise the accuracy and reliability of the weak lensing shear measurement, selecting galaxies with \((\text{S/N} > 5)\), a faint magnitude cut \(m_{\text{F606}} < 27.5\), and galaxy size \((r_h > 3\) pixels\). Our resulting weak lensing catalogue includes over 60000 objects, or roughly 65 galaxies per square arcmin. The average galaxy magnitude of this sample is \((m_{\text{F606}}) = 25.7\), implying a median redshift \(z_m \simeq 1.4\). Assuming a redshift distribution given by \(n(z) \propto z^2 \exp(-z^{1.5})\) (Baugh & Efstathiou 1993), we estimate a \(\sim 3\%\) contamination of our source galaxy catalogue from objects that are foreground to the cluster. The dilution of the signal by foreground galaxies is therefore well within the statistical noise of our analysis.

To calculate the model-free mass estimate in Equation 4 we place the background source galaxy sample at one redshift \(z_s\) taken to be the median redshift \(z_m \simeq 1.4\) of the sources. Note that for \(z_s > 1\) and \(z_l = 0.16\), the redshift of A901/2, the important distance ratio \(D_s/D_{ls}\) in Equation 1 is fairly insensitive to the value of \(z_s\). For example, a large increase of \(z_s\) from \(z_s = 1\) to \(z_s = 1.5\) increases \(D_s/D_{ls}\) by only \(\sim 6\%\). Hence for this deep analysis, where the majority of sources have redshifts \(z_s > 1\), placing all lensed galaxies at one redshift is a good approximation.

4 RESULTS

In this section we present the results of our weak lensing analysis of the A901/2 supercluster including cluster mass estimates, a comparison of three different weak lensing dark matter reconstructions and a comparison of the resulting dark matter distribution to the distribution of light in the supercluster.
4.1 Mass estimates for NFW profiles

We use spherical NFW haloes to model the weak lensing shear measured in the A901/2 field. We test two different models using the method described in section 2.3. The ‘one halo’ model centres a single NFW halo on the brightest cluster galaxy (BCG)4 in each cluster at \( z = 0.16 \). The ‘two halo’ model places a halo at the A901a BCG and the location of the infalling X-ray group A901α, a halo at the A902 BCG and at the background cluster CB1 BCG, and two halos in the SW group, SWa and SWb. There is no motivation to fit the A901b cluster with two halos and it is therefore only listed in the ‘one halo’ model upper section of the Table. The NFW ‘virial’ mass \( M_{\text{NFW}}(\rho_{\text{NFW}}) \) corresponds to a ‘virial’ radius \( r_{\text{NFW}}(\rho_{\text{NFW}}) \) which has an observed angular scale \( \theta_{\text{NFW}} \) (arcmin). For comparison with the 1 arcmin aperture model-free mass estimates \( M_{\text{ap}} \) in Table 2, \( M(\theta < 1') \) is the mass of the NFW halo enclosed by a 1 arcmin aperture, centred on (RA,Dec). The reduced \( \chi^2 \) of the fit is given in the final column.

![Table 1. Mass measurements for the A901/2 supercluster assuming the NFW spherical halo model. The ‘one halo’ model places a single NFW halo at position (RA,Dec) centred on the BCG in each cluster. The ‘two halo’ model places a halo at the A901a BCG and the location of the infalling X-ray group A901α, a halo at the A902 BCG and at the background cluster CB1 BCG, and two halos in the SW group, SWa and SWb. There is no motivation to fit the A901b cluster with two halos and it is therefore only listed in the ‘one halo’ model upper section of the Table. The NFW ‘virial’ mass \( M_{\text{NFW}}(\rho_{\text{NFW}}) \) corresponds to a ‘virial’ radius \( r_{\text{NFW}}(\rho_{\text{NFW}}) \) which has an observed angular scale \( \theta_{\text{NFW}} \) (arcmin). For comparison with the 1 arcmin aperture model-free mass estimates \( M_{\text{ap}} \) in Table 2, \( M(\theta < 1') \) is the mass of the NFW halo enclosed by a 1 arcmin aperture, centred on (RA,Dec). The reduced \( \chi^2 \) of the fit is given in the final column.](https://www.nature.com/articles/mnras32727)

<table>
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<tr>
<th>Structure</th>
<th>RA (deg)</th>
<th>Dec (deg)</th>
<th>( M_{\text{NFW}} ) (( h^{-1}10^{13}M_\odot ))</th>
<th>( r_{\text{NFW}} ) (( h^{-1}\text{kpc} ))</th>
<th>( \theta_{\text{NFW}} ) (arcmin)</th>
<th>( M(\theta &lt; 1') ) (( h^{-1}10^{13}M_\odot ))</th>
<th>( \chi^2 )</th>
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<tr>
<td>One Halo:</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>A901a</td>
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<td>22.0±4.6.0</td>
<td>1112±4.93</td>
<td>9.6±4.8</td>
<td>2.6</td>
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<tr>
<td>A901b</td>
<td>148.9889</td>
<td>−9.9841</td>
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<td>1094±4.96</td>
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</tr>
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<td>A902</td>
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<td>6.5±4.9</td>
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<td>SW group</td>
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<tr>
<td>Two Halo:</td>
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<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
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<td>149.1099</td>
<td>−9.9561</td>
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</tr>
<tr>
<td>CB1</td>
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</tr>
<tr>
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<td>572±4.133</td>
<td>4.9±4.9</td>
<td>-</td>
<td></td>
</tr>
</tbody>
</table>

4 In this analysis we define a BCG to be the brightest cluster galaxy within the peak of the clusters galaxy distribution.
distortions, a significant detection of a curl or B-mode signal indicates that ellipticity correlations exist from residual systematics. Comparing the B-mode map with the contours from the dark matter map therefore allows one to assess the reliability of each detected structure. For the maps shown in Figure 3, a 3σ detection has $\kappa \approx 0.07$, although the true significance of any peak in the distribution has to be determined by comparison to the statistics of a random Gaussian field (Van Waerbeke 2000). For a field this size, with the same number of galaxies, ellipticity distribution and smoothing scale, smoothed Gaussian noise would produce $2 \pm 3$ random $>3\sigma$ E and B mode peaks, which we discuss further in section 4.4. The most significant $>3.5\sigma$ B-mode peaks can all be linked to regions where the simple semi-time-dependent PSF modeling used in this analysis fails, as discussed in section 3.1. As an example tile 46, between A901a and A901b, was observed in two-gyro mode, out of sequence with the three-gyro main data set and three days before the other six two-gyro ACS images shown in Figure 1. The B-mode map from this ACS image exhibits the strongest $4.3\sigma$ peak.

For comparison with previous analyses, the upper right panel of Figure 3 shows a Kaiser & Squires (1993, KS93) reconstruction. The main difference seen between our preferred maximum likelihood reconstruction (upper left panel) and the KS93 reconstruction (upper right panel) is the strength of the peaks. In the

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**Figure 3.** The dark matter reconstruction of the A901/2 supercluster. The maximum likelihood dark matter map is shown in the upper left panel clearly revealing the four main superclusters structures: A901a, A901b, A902 and the SW group. This reconstruction can be compared to a Kaiser & Squires (1993, KS93) reconstruction (upper right panel), the associated B-mode or systematics reconstruction (lower left panel) and the Gray et al. (2002) ground-based dark matter reconstruction (lower right panel). The contours enclose the $-4\sigma$, $-2\sigma$, $2\sigma$, $4\sigma$, $6\sigma$, and $8\sigma$ detection regions (solid) of the STAGES maximum likelihood reconstruction (upper right and repeated on the lower panels), and the KS93 reconstruction (upper left panel). The scale bar shows the amplitude of the mapped convergence $\kappa$ which has been smoothed on 0.75 arcminute scales corresponding to $\sim130$ kpc at the supercluster redshift $z = 0.16$. The sharp edges of the maximum likelihood dark matter map and the corresponding Bmode (left hand panels) result from the data mask that is shown in the tiling pattern of Figure 1.
close-up of the four main structures of the A901/2 supercluster. This Figure compares the distribution of dark matter (shown contoured) to the luminosity weighted distribution of old red sequence galaxies defined in Wolf et al. (2005). The locations of the BCGs are shown with filled diamonds. For A901a and A901b, the maximal peak in the dark matter distribution is practically co-incident with the location of the BCG, (within 0.25 arcmin). For A902 we find two peaks in the dark matter distribution matching the two BCGs. The dark matter peaks are slightly offset from the BCGs (0.5 and 1 arcmin) due to the presence of CB1, the background cluster at a redshift of $z = 0.46$ whose location is shown in the A902 lower left panel of Figure 5 with a star. The NFW ‘two halo’ A902 and CB1 model detailed in section 4.1 predicts a shift in the observed A902 dark matter peak by $\sim 0.3$ arcmin which is consistent with what we find in the dark matter map.

For the SW group, there is again good agreement with the position of the peak in the mass distribution and the BCG, although for this group there are three local maxima in the dark matter distribution. Interestingly there are two distinct groups in the galaxy population of the SW group. There is an old red galaxy population that surrounds the BCG, as shown in the lower right panel of Figure 5. In addition there is a dusty red galaxy population (described by Wolf et al. 2005, but not shown in the Figure) that exists to the east of the BCG and coincides with the eastern dark matter peak (denoted SWa in Table 1). A more detailed analysis of the interesting relationship between the dark matter environment and the different galaxy populations will be presented in a future paper. The lower right panel of Figure 5 also shows one case of a significant density of red galaxies without a peak in the dark matter distribution. Towards the edge of the STAGES imaging, the noise in our dark matter map grows rapidly, and at the location of this galaxy group the noise is twice the noise level at the SW group. As this group is likely to be less massive than the SW group, which is detected at 5$\sigma$, we are not surprised that this group is undetected in our dark matter map.

The A901a upper left panel of Figure 5 shows a significant...
A comparison of mass and light in the main structures of the A901/2 supercluster; A901a (upper left), A901b (upper right), A902 (lower left) and the SW group (lower right). The locations of the brightest cluster galaxies are shown (filled diamonds), in addition to the location of the in-falling X-ray group A901α (filled triangle), and the location of the higher redshift z = 0.46 cluster CB1 (star). Local maxima in the dark matter map are shown with a cross.

Comparing the local maxima in the A901b and A902 distribution with the light maps we find that the substructure in the dark matter maps are often associated with substructures in the galaxy distribution. The only striking discrepancy is a luminous peak to the north west of A902, seen in Figure 4. This luminous peak results from a single, very luminous dusty red galaxy that is brighter than the BCG and is likely to be infalling on A902.

In Table 2 we list mass and mass-to-light ratios for the main structures shown in Figure 5. As discussed in section 4.1, these structures are far from the spherically symmetric NFW models that are often used to constrain models. We therefore use a model-free mass estimate given by Equation 4, defining the enclosed region using the 1σ and 3σ contours shown in Figure 5. For comparison with the ground-based analysis of Gray et al. (2002) and the NFW analysis of section 4.1 we also list the mass enclosed by a 1 arcmin circular aperture (denoted ‘ap’) centered on each clusters BCG. To estimate the contribution of systematic error to our mass estimate we follow the conservative prescription that is often used in the analysis of weak lensing by large-scale structure (see for example Benjamin et al. 2007), calculating errors by adding the random error (listed as the first mass error in Table 2) in quadrature with the B-mode signal, shown in the lower left panel of Figure 3 and listed as the second mass error in Table 2. The systematic error dominates the random error in this analysis.

We find A901a and A901b to be the most massive systems in the supercluster with masses \( \sim 6.5 h^{-1}10^{13} M_\odot \) and mass-to-light ratios \( \sim 150 \) for the full extended 1σ region. A902 and the SW group have similar masses, roughly half the mass of the A901 pair at \( \sim 3.5 h^{-1}10^{13} M_\odot \). We find A901b to be the most extended structure in the system, and the SW group is the most compact.

The mass-to-stellar mass ratios \( M/M_* \) of each structure are given in the final column of Table 2. These mass ratios were calculated with a Hubble parameter \( h = 0.7 \), assuming a Kroupa et al. (1993) initial mass function. The results are equivalent to within 10% of the same result derived using a Chabrier (2003) or a Kroupa (2001) initial mass function. We find mass-to-stellar mass ratios \( M/M_* \) that are similar to the ratios found for massive elliptical galaxies at this redshift (Hoekstra et al. 2005; Mandelbaum et al. 2006; Heymans et al. 2006), although a direct comparison is hard to draw as the results from the massive elliptical galaxies measure NFW virial mass to stellar mass ratios instead of the model free mass ratio estimates that we present here.

Figure 6 shows the variation of the mass-to-stellar mass ratio across each of the main structures in A901/2 on a log scale, compared to the mass distribution (shown contoured). Note that negative regions in the mass reconstruction have been set to zero in this figure. Moving out from the central BCG (shown contoured), we find that the mass-to-stellar mass ratio initially increases, as the stellar mass decreases more rapidly than the halo mass. This can be compared to the light map in Figure 5. Continuing out further, the mass-to-stellar mass ratio then rapidly decreases as the dark matter mass decreases. This Figure shows some regions of very high mass-to-stellar mass ratio regions (log \( M/M_* > 2 \)), but the reader should note the significance of the mass detected in these regions (shown contoured) which is less than 2σ in all cases.

Comparing our convergence \( \kappa \) mass reconstruction results to our NFW shear analysis, we find very good agreement between...
the mass measured within 1 arcmin of each clusters BCG. This provides an important verification of our dark matter reconstruction method. In comparing our results to the previous ground-based lensing analysis of the A901/2 supercluster we must first consider the assumption made by Gray et al. (2002) that \( \kappa \ll 1 \) and hence \( \gamma = g \), where \( g \) is the measured reduced shear given above Equation 2. For the main structures in A901/2 this would result in an overestimate of cluster mass by \( \sim 15\% \). Taking this overestimate into account, our mass estimates are consistent with Gray et al. (2002) as can be seen from the space/group mass map comparison in the lower right panel of Figure 3. The mass-to-light ratio measurements for A901b and A902 disagree however at the 3σ level. The difference arises from dramatic improvements in the selection of the cluster galaxies from the COMBO-17 data, in comparison to the previous two-band optical cluster selection of Gray et al. (2002). This improvement removes the striking difference between the cluster mass-to-light ratio measurements found by Gray et al. (2002). Our results show a mass-to-light ratio within an aperture of 1 arcmin of \( M(\leq 1') \sim 100h^{-1}M_{\odot}L_{\odot} \) to be a good description for all the main structures in the supercluster. We find a similar result for the mass-to-stellar mass ratio where \( M/M_\ast(\leq 1') \sim 25 \) for all the main structures in the supercluster.

Table 2. Model-free mass measurements for the main structures in the A901/2 supercluster. The enclosed masses \( M \), mass-to-light ratios \( M/L \) and mass-to-stellar mass ratios \( M/M_\ast \) are given for the regions defined by the 1σ and 3σ contours in the signal-to-noise maps. The area of these regions is given in the second column (square arcmin). For comparison with the ground-based analysis of Gray et al. (2002) we also list the mass and mass-to-light enclosed by a 1 arcmin circular aperture (denoted ‘ap’) centered on the cluster BCG. The quoted mass errors are listed with the random noise error followed by a conservative estimate of the systematic error. The mass-to-light ratio error and mass-to-stellar mass ratio error include only the error on the mass with the random and systematic parts added in quadrature.

<table>
<thead>
<tr>
<th>Region</th>
<th>( M )</th>
<th>( M/L )</th>
<th>( M/M_\ast )</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>( h^{-1}10^{13}M_{\odot} )</td>
<td>( h^{-1}M_{\odot}/L_{\odot} )</td>
<td>( h = 0.7 )</td>
</tr>
<tr>
<td>A901a</td>
<td>6.34 ± 0.06 ± 1.25</td>
<td>134.5 ± 26.5</td>
<td>32.9 ± 6.5</td>
</tr>
<tr>
<td>1σ, 16.9</td>
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<td></td>
<td></td>
</tr>
<tr>
<td>3σ, 8.2</td>
<td>4.28 ± 0.04 ± 0.68</td>
<td>135.2 ± 21.6</td>
<td>31.4 ± 5.0</td>
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<tr>
<td>ap, 3.1</td>
<td>2.00 ± 0.02 ± 0.20</td>
<td>98.6 ± 10.2</td>
<td>21.7 ± 2.2</td>
</tr>
<tr>
<td>A901b</td>
<td>6.75 ± 0.07 ± 2.51</td>
<td>163.0 ± 60.7</td>
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</tr>
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<td>1σ, 21.7</td>
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<td></td>
</tr>
<tr>
<td>3σ, 6.5</td>
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<td>206.3 ± 82.3</td>
<td>50.8 ± 20.3</td>
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<td>ap, 3.1</td>
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<td>148.7 ± 63.2</td>
<td>34.7 ± 14.8</td>
</tr>
<tr>
<td>A902</td>
<td>3.40 ± 0.05 ± 0.51</td>
<td>108.2 ± 16.4</td>
<td>27.6 ± 4.2</td>
</tr>
<tr>
<td>1σ, 12.6</td>
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<td></td>
</tr>
<tr>
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<tr>
<td>SW group</td>
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<td>126.2 ± 37.4</td>
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</table>

Table 3. Peak statistics; comparing the number of peaks as a function of significance in a smoothed random Gaussian field (Noise), the reconstructed dark matter map (Signal) and the reconstructed B-mode map. We use the peak profile to further distinguish between a noise peak and a true peak. Signal (p) and B-mode (p) list the number of peaks in our substructure sample, where the peak has less than 33% chance of being a noise peak. The Signal and B-mode peaks are very rarely co-incident, but as a conservative measure we define a high confidence sample, (Signal (b)), where the B-mode at the peak must be less than half the amplitude of the signal.

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<th>( \kappa_{\text{peak}} &lt; )</th>
<th>-6σ</th>
<th>-4σ</th>
<th>-2σ</th>
<th>2σ</th>
<th>4σ</th>
<th>6σ</th>
<th>8σ</th>
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<td>19</td>
<td>6</td>
<td>0</td>
<td>0</td>
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<tr>
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<td>39</td>
<td>11</td>
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<td>1</td>
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<tr>
<td>Bmode (p)</td>
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<td>36</td>
<td>45</td>
<td>2</td>
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<td>33</td>
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<td>1</td>
</tr>
<tr>
<td>Bmode (b)</td>
<td>0</td>
<td>0</td>
<td>6</td>
<td>33</td>
<td>2</td>
<td>0</td>
<td>0</td>
</tr>
</tbody>
</table>

4.4 Supercluster substructure

In this section we investigate the lower significance peaks in the dark matter distribution that are not associated with the cores of the supercluster structures discussed above. Table 3 lists the number of local maxima and minima in the dark matter reconstruction for different significance levels and compares them to what we find in our B-mode reconstruction and what we would expect from a smoothed random Gaussian field using Equation (41) from Van Waerbeke (2000). The high significance peaks \( \kappa_{\text{peak}} > 3\sigma \) are all associated with the cores of the four main supercluster structures. However we can see that we have a significant number of \( \kappa_{\text{peak}} > 2\sigma \) peaks that cannot be explained by random noise alone. There are a comparable number of \( \kappa_{\text{peak}} > 2\sigma \) peaks in the B-mode map, but comparing the location of E and B mode peaks allows one to assess the reliability of the lower significance E mode detections.

In order to distinguish between noise peaks and true peaks in the mass distribution, it is useful to add morphological information about the profile of the peaks. The mean profile and dispersion of a noise peak is given by Equation (47) of Van Waerbeke (2000). Comparing the measured profile around each detected peak with the mean noise profile allows for the calculation of the probability that a peak with a given significance and shape is a noise fluctuation, (using Equation (45) of Van Waerbeke (2000)). In Table 3 we list the number of peaks that have a less than 33% probability of being a random noise fluctuation. These are consistent with the numbers expected from comparing the number of detected peaks with the number expected from noise.

To define a low-significance \( 2\sigma < \kappa_{\text{peak}} < 4\sigma \) substructure sample we use high confidence selection criteria where the peak must have less than 33% probability of being a random noise fluctuation, and the B-mode at the location of the peak must be less than half the amplitude of the E-mode. The last row of Table 3 lists the number of peaks that meet these criteria for different significance levels, leaving 14 ‘substructure’ peaks with \( 2\sigma < \kappa_{\text{peak}} < 4\sigma \). Note that the 11 peaks with \( \kappa_{\text{peak}} > 4\sigma \) are all associated with the central regions of the four main structures in the supercluster.
discussed in section 4, as shown by the marked crosses in Figure 5 that are enclosed by the 4σ contour.

Assuming all 14 substructure peaks in the dark matter map are associated with the supercluster, we can calculate a mass for these halos using Equation 4. We find an average mass of $M(<0.75') = 0.29 \pm 0.02 h^{-1}10^{13}M_\odot$ for the $2\sigma < \kappa_{\text{peak}} < 3\sigma$ group, and $M(<0.75') = 0.46 \pm 0.06 h^{-1}10^{13}M_\odot$ for the $3\sigma < \kappa_{\text{peak}} < 4\sigma$ group. Table 4 lists the number of peaks that are associated with cluster galaxies where $L(<0.75') > 10^{7}L_\odot$. We find that half of these peaks are associated with galaxies in the dark matter map that are detected at $>4\sigma$. The remaining peaks in the distribution are likely to be associated with galaxy groups at higher redshift (Falter et al. 2007), supported by the discovery of several co-incident groups of galaxies at $z \sim 0.45$.

The dark matter map presented in this paper will form the basis of future studies of galaxy morphology and galaxy type in an over-dense dark matter environment. The high resolution HST data set is also ideal for the analysis of weak gravitational flexion, a third order weak lensing effect that will be very effective at probing the sub-structures that were resolved in this weak shear analysis (see for example Bacon et al. 2006). Comparing the results of this analysis with the previous ground-based analysis clearly demonstrates the importance of space-based observations for future weak lensing studies.

5 DISCUSSION AND CONCLUSION

From a weak lensing analysis of deep Hubble Space Telescope data, we have reconstructed a high-resolution map of the dark matter distribution in the Abell A901/A902 supercluster. We find that the maximal peaks in the dark matter distribution are very well matched to the locations of the brightest cluster galaxies in the most massive structures in the supercluster. These structures are A901a, A901b, A902 and the South West group, all of which are detected in our dark matter map at high significance.

Owing to the high number density of resolved objects in the HST data, we have been able to increase the resolution of the dark matter map by a factor of $\sim 2.5$ over the ground-based analysis of Gray et al. (2002). This has allowed us to resolve the morphology of the dark matter structures, finding profiles that are far from the spherically symmetric NFW models that are typically used to model such systems. We find local maxima in the dark matter distribution around the main structures, that are also seen in the distribution of galaxies. Furthermore we see a significant extension in the dark matter distribution around A901a, in the direction of an in-falling X-ray group called A901α (Gray et al. 2007a).

We have presented mass, mass-to-light and mass-to-stellar mass ratio estimates for each of the main structures, finding A901a and A901b to be the most massive clusters in the system with $M(<1') \sim 2 h^{-1}10^{13}M_\odot$. Contrary to the analysis of Gray et al. (2002) we find no evidence for the variation of the mass-to-light ratio or the mass-to-stellar mass ratio between the different clusters measured on the central 1 arcmin ($\sim 170$ kpc) of each cluster. We have shown the variation of the mass-to-stellar mass ratio across the clusters, finding an initial rise in $M/M_\odot$ as a function of distance from the clusters central BCG, followed by a steep decrease.

We have investigated the less significant substructures in the dark matter map that are detected at $<4\sigma$. The profile of these peaks with what is expected from a random noise peak we have selected a sample of substructures where the likelihood of those peaks being noise or a result of an imperfect PSF correction is low. We find that half of these peaks are associated with galaxies in the cluster, yielding mass-to-light ratios that are comparable to the mass-to-light ratios found in the main structures in the supercluster. The remaining peaks in the distribution are likely to be associated with galaxy groups at higher redshift (Falter et al. 2007), supported by the discovery of several co-incident groups of galaxies at $z \sim 0.45$.

We have presented mass, mass-to-light and mass-to-stellar mass ratio estimates for each of the main structures, finding A901a and A901b to be the most massive clusters in the system with $M(<1') \sim 2 h^{-1}10^{13}M_\odot$. Contrary to the analysis of Gray et al. (2002) we find no evidence for the variation of the mass-to-light ratio or the mass-to-stellar mass ratio between the different clusters measured on the central 1 arcmin ($\sim 170$ kpc) of each cluster. We have shown the variation of the mass-to-stellar mass ratio across the clusters, finding an initial rise in $M/M_\odot$ as a function of distance from the clusters central BCG, followed by a steep decrease.

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